

## Experimental Study of a Small-cell Drift Chamber Operated with Helium-Based Gases

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**Abstract** The performance of a small-cell drift chamber operated with two helium-based gas mixtures, He/C<sub>3</sub>H<sub>8</sub>(60/40) and He/CH<sub>4</sub>(60/40), was studied using cosmic rays, respectively. The gas mixture Ar/CO<sub>2</sub>/CH<sub>4</sub>(89/10/1) was also tested for comparison. The spatial resolution and  $dE/dx$  resolution of the helium-based gases are better than those of Ar/CO<sub>2</sub>/CH<sub>4</sub>(89/10/1). A good spatial resolution of 110 $\mu$ m and a  $dE/dx$  resolution of 6%—7% with 30—40 samples were obtained with He/C<sub>3</sub>H<sub>8</sub>(60/40).

**Key words** drift chamber, small cell, helium-based gas mixture

### 1 Introduction

The BES III<sup>2)</sup> experiment is designed to perform high-precision measurements and new physics search at the tau-charm factory BEPC II<sup>3)</sup>. This places stringent demands on the performance of the drift chamber for BES III. It must provide high efficiency, good momentum resolution and excellent adaptability for high count rate. In addition, the drift chamber is required to serve particle identification with ionization loss( $dE/dx$ ) measurements. The solution to meet the above requirements is a small-cell drift chamber made of low mass material<sup>2)</sup>. The small cell design has many advantages with good spatial resolution, good  $dE/dx$  resolution, long aging lifetime and fast response. The latter two are essential for the operation under the high luminosity condition of BEPC II. The majority of charged particles produced in BEPC has the momentum lower than 1GeV/ $c$ . The predominant limitation on the momentum resolution is the error due to the multiple scattering in the drift chamber. It can be minimized by reduc-

ing the amount of material in the drift chamber. So the light gas is preferred. The helium-based gas is the good candidate for such kind of gas. In addition, the helium-based gas can reduce the background from large amount of synchrotron radiation produced in the collider with high luminosity for its small photo-electric cross section.

So far, many kinds of helium-based gases have been studied<sup>[1-5]</sup>. The small-cell drift chamber has also been successfully used in many large experiments of high energy physics<sup>[6-8]</sup>.

In order to study the performance of the small-cell drift chamber filled with helium-based gas mixtures, we built a prototype drift chamber with the cell size similar to the design of the BES III drift chamber and carried out some tests using cosmic rays. In this paper, we present the test results of the prototype drift chamber with two helium-based gas mixtures: He/C<sub>3</sub>H<sub>8</sub>(60/40) and He/CH<sub>4</sub>(60/40), using cosmic rays. The gas mixture of He/C<sub>3</sub>H<sub>8</sub>(60/40) is the perfect candidate for the BES III drift chamber with its excellent properties of good spatial and  $dE/dx$  resolution and long radiation length

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2)BES III Preliminary Design Report. IHEP-BEPC II -SB-13, 2004.

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(550m). The gas mixture of He/CH<sub>4</sub> (60/40) is also interesting for its low price, easy access and very long radiation length(800m). In addition, the properties of Ar/CO<sub>2</sub>/CH<sub>4</sub> (89/10/1), which was used by the BES II drift chamber, were measured for comparison.

## 2 Experimental setup

The test chamber consists of 14 sense wire layers with 5 cells per layer, as shown in Fig.1. The cells are all in standard square shape to simplify the drift distance-time calibration, with the size of 16.4mm × 16.4mm. Each cell has one sense wire surrounded by a square grid of 8 field wires which

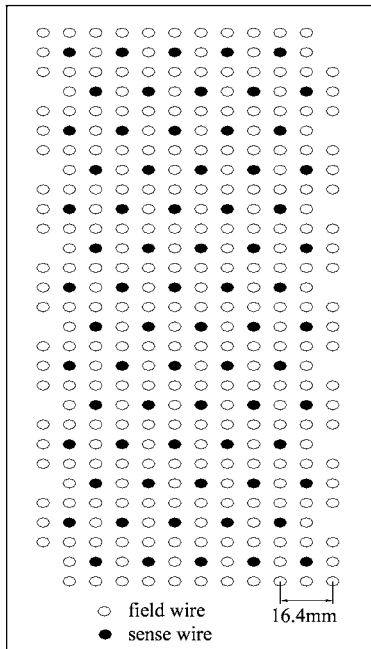


Fig. 1. The cell configuration of the test chamber.

are shared by adjacent cells. The sense wire is made of gold-plated tungsten(3% Re), 25 $\mu$ m in diameter, tensed with a weight of 20g. The field wire is made of gold-plated aluminum with a diameter of 110 $\mu$ m. A tension of 60g is applied to the field wires to match the gravitational sag of the sense wires. The wires are about 60cm long. All layers are axial with half cell staggering for left-right ambiguity resolution. The end-plates and the walls of the test chamber are all made of aluminum. A piece of 1 cm long slit is machined in the middle of the upper wall to allow the use of <sup>55</sup>Fe X-ray to test the operation of the chamber. The high voltage is applied to the sense wires by a printed board at one end-plate, while the

preamplifier boards are mounted on the other end-plate. The field wires are all connected to ground directly. The operating voltages for various gases are determined in such a way that the gas gains are similar.

The chamber was tested with various gases using cosmic rays. The test setup and electronics have been described in detail in Ref. [9].

The gas mixture He/C<sub>3</sub>H<sub>8</sub> (60/40) was supplied by mass flowmeters. For He/CH<sub>4</sub> (60/40) and Ar/CO<sub>2</sub>/CH<sub>4</sub> (89/10/1), pre-mixed gases were used.

## 3 Results and discussion

### 3.1 Drift distance-time relation

The data of cells used in analysis are required to have a drift time less than 500ns and a charge larger than the probable pedestal by 5 times' standard deviations. These cells are called hit cells. The probable pedestal and standard deviation are obtained by fitting the pedestal distribution to a Gaussian.

The electric field in the square small cell is not uniform because of the cell's geometric structure. As a result, the drift distance-time relation ( $x-t$  relation) for the cell is non-linear. So an iterative algorithm is used to determine the  $x-t$  relation. With a rough  $x-t$  relation and after simple pattern recognition the track is fitted to the circles centered at the sense wires and with the radii corresponding to the measured drift distance of hit cells. Then the residuals for each drift time bin, defined as the signed difference between the drift distance calculated with drift time using  $x-t$  relation and the distance from the sense wire to the fitted track, are histogrammed. Each cell used for the calculation of residuals is excluded from the track fit to make the residual distribution bias free. The  $x-t$  relation is then corrected in each time bin by the average residual value. The new  $x-t$  relation is used in the next iteration. The  $x-t$  relation converges after a few iterations. In the iteration process, the adjustments of fine TO and wire position are performed channel by channel, based on the residual distributions on the left and right side of sense wires to make the distributions center at about zero.

The  $x-t$  relations for various gas mixtures are all nonlinear as shown in Fig.2. He/CH<sub>4</sub> (60/40) has a smaller non-linearity compared with He/C<sub>3</sub>H<sub>8</sub> (60/40). This is consistent with the fact that He/CH<sub>4</sub> (60/40) has a lower electric field for the saturation of the electron drift velocity. The drift ve-

locities of helium-based gases are slower than Ar/CO<sub>2</sub>/CH<sub>4</sub> (89/10/1) but still fast enough so that the maximum drift time for about 8mm distance at normal incident angle is less than 300ns. It is acceptable for BEPC II .

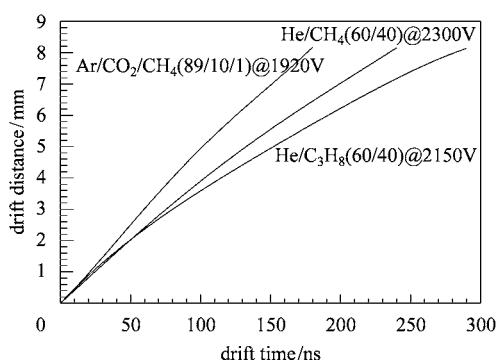


Fig. 2.  $x-t$  relations for various gases.

### 3.2 Spatial resolution

The spatial resolution is determined from the residual distribution. The width of the residual distribution has two contributions which add in quadrature: the intrinsic spatial resolution and the error of the track fit. So the residual has to be corrected to get the intrinsic spatial resolution without the track fit error. The time walk correction is also applied to improve the spatial resolution, based on the correlation of the timing and the pulse height.

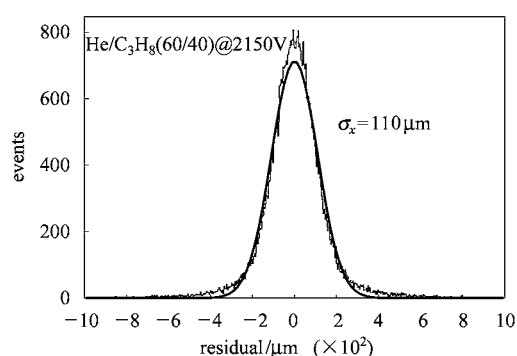


Fig. 3. The typical residual distribution.

The typical overall residual distribution after the above two corrections is shown in Fig. 3. The average spatial resolution is extracted by fitting the distribution to a Gaussian as shown in Fig. 3. The distribution has a pronounced non-Gaussian tail. It originates from the large spread of the drift time for the electrons ionized along the incident track to reach

the sense wire, which is called non-isochronous behavior. This is the primary factor limiting the spatial resolution for the small cell.

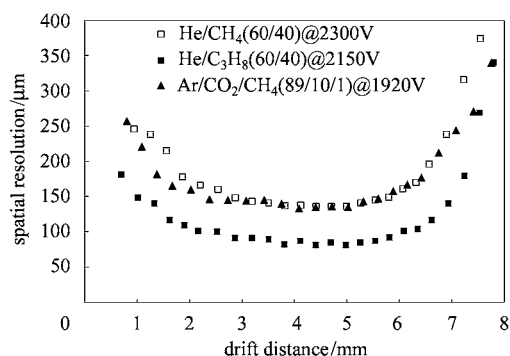


Fig. 4. The spatial resolution as a function of drift distance.

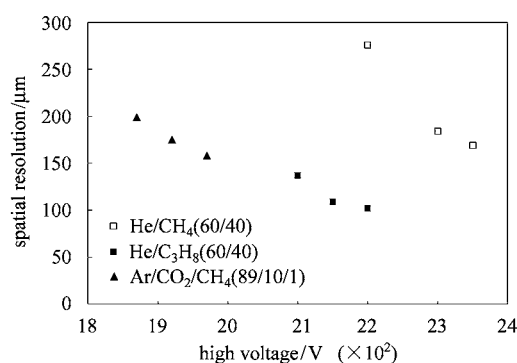


Fig. 5. The spatial resolutions of various gases as a function of the operating voltage.

The dependence of the spatial resolution on the drift distance is shown in Fig. 4. The spatial resolution deteriorates near the sense wire because of the statistic fluctuation of primary ionization and the sharp degradation near the cell edge is due to the distortion of the electric field. The spatial resolution for various gases is summarized in Fig. 5 as a function of the operating voltage. It is universal that the spatial resolution is improved significantly with the increase of the operating voltage when the high voltage is relatively low. He/C<sub>3</sub>H<sub>8</sub> (60/40) achieves a spatial resolution of 110 $\mu$ m at 2150V. The spatial resolution of He/C<sub>3</sub>H<sub>8</sub> (60/40) is much better than that of He/CH<sub>4</sub> (60/40). This is due to the fact that He/C<sub>3</sub>H<sub>8</sub> (60/40) provides a much larger number of primary ions (30ions/cm) than that of He/CH<sub>4</sub> (60/40) (10ions/cm). Although Ar/CO<sub>2</sub>/CH<sub>4</sub> (89/10/1) has large primary ions (30ions/cm), its spatial resolution is very bad compared

with He/C<sub>3</sub>H<sub>8</sub>(60/40). The reason may be that the operating voltage is low and the fluctuation of the drift velocity in the cell is big. The incident azimuthal angle dependence of the spatial resolution was also studied. The result shows no significant deterioration up to 13°.

### 3.3 dE/dx resolution

For dE/dx measurement, the track was also required to pass through about 25cm of lead. This eliminated the tracks with low momenta which have a large range of dE/dx. Before calculating the dE/dx resolution, some corrections are applied on the pulse height. At first the path length correction is done using the reconstructed track. Then the gas gain variation due to the pressure or temperature change is corrected run by run. The relative gain difference between sense wires is also corrected channel by channel.

In order to obtain the dE/dx resolution with the samples similar to the expectation from the BES III drift chamber (about 40), several events are combined for the measurement. The dE/dx resolution is measured using the truncation technique<sup>[10]</sup> to minimize the contribution of the well-known Landau tail<sup>[11]</sup> to dE/dx resolution. With the truncation technique, a Gaussian-like dE/dx spectrum is produced. The spectrum is then fitted to a Gaussian. The dE/dx resolution is defined as the ratio of the standard derivation of the Gaussian to the peak value.

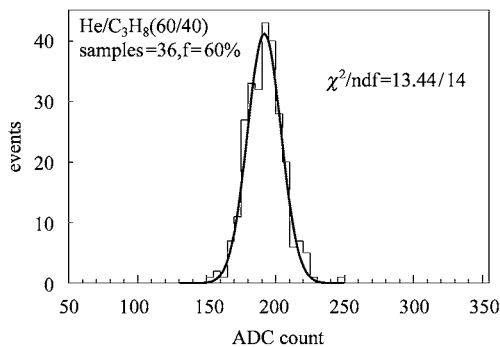


Fig. 6. The 60% truncated mean pulse height distribution.

Fig. 6 shows the truncated mean pulse height distribution with 36 samples, 60% truncation fraction, for He/C<sub>3</sub>H<sub>8</sub>(60/40) at 2150V. The distribution is quite Gaussian-like. The dE/dx resolution for various gases at various operating voltages is shown in Fig. 7. He/C<sub>3</sub>H<sub>8</sub>(60/40) has the best resolution for the presence of a large fraction of hydrocarbon. Despite the less primary ionization number, He/CH<sub>4</sub>(60/40)

has a better dE/dx resolution than Ar/CO<sub>2</sub>/CH<sub>4</sub>(89/10/1). This is because hydrocarbon gases have a very good dE/dx resolution<sup>[12]</sup>. The dependence of dE/dx resolution on the samples is also studied for He/C<sub>3</sub>H<sub>8</sub>(60/40) at 2150V, as shown in Fig. 8. A resolution of 6%—7% can be obtained with 30—40 samples.

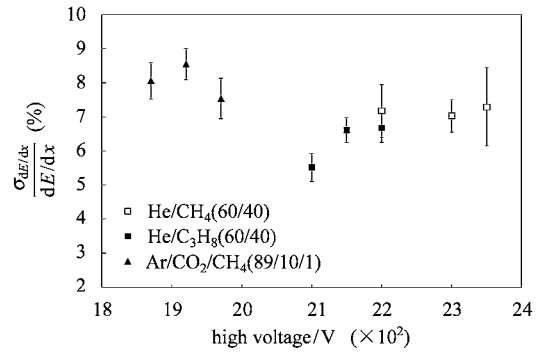


Fig. 7. The dE/dx resolution of various gases at several operating voltages.

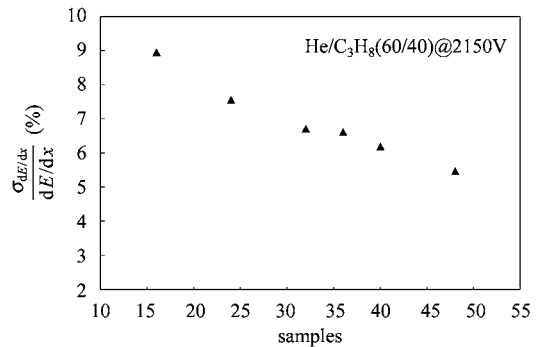


Fig. 8. The dE/dx resolution as a function of samples for He/C<sub>3</sub>H<sub>8</sub>(60/40).

### 3.4 Cell efficiency

The cell efficiency is calculated to be the ratio of the tracks with hits in the cell within a 5-σ window to the total tracks passing through the cell, where σ is the average spatial resolution. Fig. 9 shows the cell efficiency as a function of the distance from the sense wire. The efficiency drops very quickly near the cell boundary. The drop arises from the small and slow signal when the track passes the cell boundary due to the very dispersed drift time for primary ionized electrons. This effect will be enhanced in the presence of the magnetic field. The efficiency for various gases at various operating voltages is shown in Fig. 10. The efficiency gets

higher with the increase of the operating voltage and then arrives at a plateau. The drift chamber should operate at the plateau of efficiency. He/C<sub>3</sub>H<sub>8</sub>(60/40) achieves a better efficiency than the other gases with the similar gas gain.

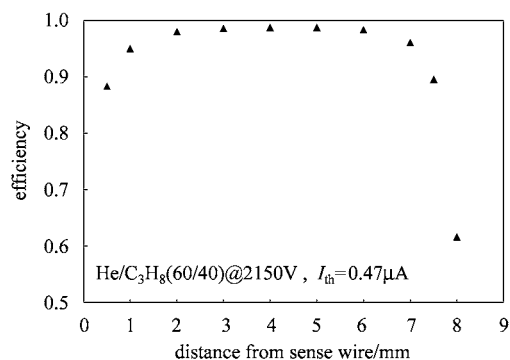


Fig. 9. The cell efficiency as a function of the distance from the sense wire.

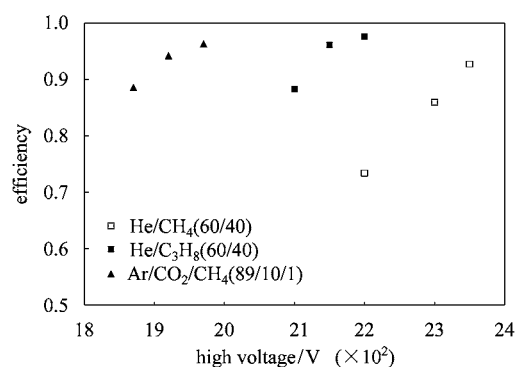


Fig. 10. The cell efficiency for various gases as a function of the operating voltage.

### 3.5 Cross-talk

The cross-talk between the neighboring cells is studied by comparing the ADC pedestal of one cell in the case of no hit in its neighboring cells with the case of having hits in the neighboring cells. A cross-talk of about  $-2.0\%$  between the neighboring cells is found. It indicates that the major contribution to the cross-talk is the signal induced on the nearby sense wire by the motion of the positive ions produced in the avalanche. As the positive ions motion relative to the nearby sense wire is opposite to the sense wire in avalanche, the induced signal has an opposite sign. Anyway the cross-talk for small cell is small and does not produce spurious hits thanks to its opposite sign.

## 4 Conclusion

We have tested the performance of a small-cell drift chamber prototype with two helium-based gas mixtures: He/C<sub>3</sub>H<sub>8</sub>(60/40) and He/CH<sub>4</sub>(60/40) as well as the mixture of Ar/CO<sub>2</sub>/CH<sub>4</sub>(89/10/1), using cosmic rays. Our study shows that the helium-based gases provide better spatial and  $dE/dx$  resolutions than Ar/CO<sub>2</sub>/CH<sub>4</sub>(89/10/1) and a good spatial resolution of  $110\mu\text{m}$  and a  $dE/dx$  resolution of  $6\% - 7\%$  with 30—40 samples can be obtained using the small cell with the helium-based gas He/C<sub>3</sub>H<sub>8</sub>(60/40). It indicates that the BESIII drift chamber can be considered to be a small-cell drift chamber operated with He/C<sub>3</sub>H<sub>8</sub>(60/40) to satisfy its performance requirements.

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## 一个使用氨基混合气的小单元漂移室的实验研究

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**摘要** 利用宇宙线分别研究了一个小单元漂移室在 He/C<sub>3</sub>H<sub>8</sub>(60/40)和 He/CH<sub>4</sub>(60/40)两种氨基混合气下的性能. 作为比较还测试了在 Ar/CO<sub>2</sub>/CH<sub>4</sub>(89/10/1)混合气下的性能. 研究表明上述氨基混合气的空间分辨和 dE/dx 分辨好于 Ar/CO<sub>2</sub>/CH<sub>4</sub>(89/10/1)混合气. 使用 He/C<sub>3</sub>H<sub>8</sub>(60/40)可以获得 110 $\mu$ m 的空间分辨, 在 30—40 次取样下 dE/dx 分辨可以达到 6%—7%.

**关键词** 漂移室 小单元 氨基混合气